

# 1 Search for correlations of high-energy neutrinos 2 and ultra-high energy cosmic rays

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[http://www.auger.org/archive/authors\\_icrc\\_2019.html](http://www.auger.org/archive/authors_icrc_2019.html)

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The sources of ultra-high energy cosmic rays (UHECRs) are still one of the main open questions in high-energy astrophysics. If UHECRs are accelerated in astrophysical sources, they are expected to produce high-energy photons and neutrinos due to the interaction with the surrounding astrophysical medium or ambient radiation. In particular, neutrinos are powerful probes for the investigation of the region of production and acceleration of UHECRs since they are not sensitive to magnetic deflections nor to interactions with the interstellar medium. The results of three different analyses that correlate the very high-energy neutrino candidates detected by IceCube and ANTARES and the highest-energy cosmic rays measured by the Pierre Auger Observatory and the Telescope Array will be discussed. The first two analyses use a sample of high-energy neutrinos from IceCube and ANTARES selected to have a significant probability to be of astrophysical origin. The first analysis cross-correlates the arrival directions of these selected neutrino events and UHECRs. The second one is a stacked likelihood analysis assuming as stacked sources the high-energy neutrino directions and looking for excesses in the UHECR data set around the directions of the neutrino candidates. The third analysis instead uses a larger sample of neutrinos selected to look for neutrino point-like sources. It consists of a likelihood method that looks for excesses in the neutrino point-source data set around the directions of the highest-energy UHECRs.

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## 3 1. Introduction

4 Galactic accelerators like supernova remnants are nowadays believed to be the most likely  
 5 sources for cosmic rays (CRs) below  $10^{15}$  eV [1]. On the other hand, ultra-high energy (above  
 6  $10^{18}$  eV) cosmic rays (UHECRs) originate from some yet-unidentified extra-galactic sources, as  
 7 indicated by the recent Pierre Auger Observatory measurement of the first statistically significant  
 8 large-scale anisotropy above 8 EeV [2]. The most promising sources, although contra-  
 9 dictory in some aspects [3, 4, 5, 6, 7, 8], are active galactic nuclei, gamma ray bursts and magnetized  
 10 and fast-spinning neutron stars. UHECRs accelerated in astrophysical sources are expected to pro-  
 11 duce high-energy photons and neutrinos when interacting with the ambient matter and radiation. In  
 12 particular, due to their tiny cross section and their insensitivity to (inter-)galactic magnetic fields,  
 13 neutrinos constitute an excellent probe to investigate the origin of UHECRs. A multi-messenger  
 14 approach might hence lead to a deep insight in the search of UHECR origin and their acceleration  
 15 mechanisms. In these proceedings, we present the results of three analyses searching for a com-  
 16 mon origin of UHECRs and high-energy neutrinos using data from IceCube Neutrino Observatory,  
 17 the ANTARES Collaboration, the Pierre Auger Observatory and the Telescope Array (TA) Col-  
 18 laboration. The analyses are (1) a cross-correlation analysis that scans angular distances between  
 19 UHECRs and high-energy neutrinos, (2) a neutrino-stacking correlation analysis with UHECR di-  
 20 rections and (3) a UHECR-stacking correlation analysis with neutrino directions.

## 21 2. Observatories and data samples

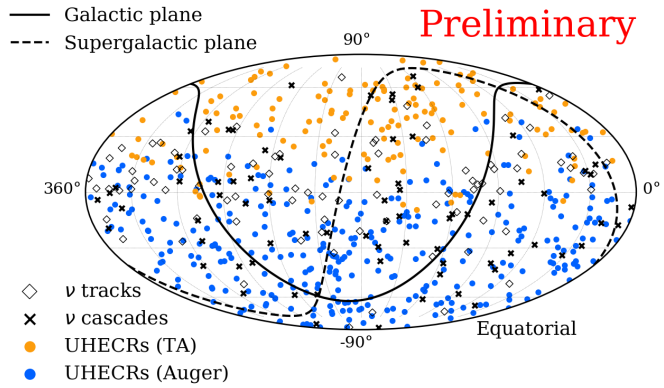
22 **IceCube** [9] is a  $1\text{-km}^3$  sized neutrino detector optimized for neutrino energies above  $\sim 100$   
 23 GeV, located at the geographic South Pole at about 1.5 to 2.5 km deep in the glacial ice. It consists  
 24 of 86 strings instrumented by 5160 photomultiplier tubes housed together with on-board digiti-  
 25 zation modules in pressure resistant spheres. The first and the second analysis, which will be  
 26 presented in the following, combine different IceCube datasets: (i) the 7.5-year data set of High-  
 27 Energy Starting Events (HESE) [10], (ii) the 9-year sample of Extremely High-Energy event alerts  
 28 (EHE) [11] and (iii) a complementary 7-year sample of through-going muons induced by charged-  
 29 current interactions of  $\nu_\mu$  candidates from the Northern sky [12]. The HESE sample is composed  
 30 by 76 shower-like events, characterized by an average angular resolution of  $\sim 15^\circ$  above 100 TeV,  
 31 and 26 track-like events, with an average angular resolution of  $\sim 1^\circ$  [13]. The IceCube realtime  
 32 neutrino alert system is based on HESE and EHE selection by analyses looking for cosmogenic  
 33 neutrinos [14]. The HESE alerts resulted in the evidence for the first neutrino emission in coinci-  
 34 dence with a high-energy gamma-ray emission from the blazar named TXS 0506+056 [15]. The  
 35 EHE analysis discovered the first observed PeV-scale neutrinos [16]. The EHE alert event selec-  
 36 tion is sensitive to energies from about 500 TeV to 10 PeV and targets track-like events, which  
 37 have good angular resolution ( $\leq 1^\circ$ ). The used EHE sample is composed by 20 track-like events.  
 38 Finally, the through-going muons are composed by 35 tracks with  $E \gtrsim 200$  TeV, corresponding to  
 39 7 years of data from the 8-year sample presented in [12]. Figure 1 shows the arrival directions of  
 40 neutrino track- and cascade-like events described above, together with the ANTARES high-energy  
 41 neutrinos and Auger and TA UHECR events described in the following. The third analysis uses (i)  
 42 the 7-year neutrino point-source sample [17] and (ii) the latest 3.5 years of the gamma-ray follow-  
 43 up (GFU) sample [18]. This combined track-like sample, selected for point-like source searches,

44 consists of 1.4 million events recorded between 2008 and 2018. These are dominated in the North-  
 45 ern hemispheres by atmospheric  $\nu_\mu$  and in the Southern hemisphere by atmospheric downgoing  
 46 muons. The angular resolution is  $< 0.5^\circ$  above TeV energies [17].

47 **ANTARES** [19] is a neutrino telescope located in the Mediterranean Sea, composed by 12 ver-  
 48 tical strings anchored at the sea floor at a depth of  $\sim 2400$  m, covering a total volume of  $\sim 0.03$  km<sup>3</sup>.  
 49 The strings are equipped with a total of 885 optical modules, each one housing a photomultiplier  
 50 tube. The events used in analyses (1) and (2) are selected from the 9-year point-source sample [20],  
 51 recorded between January 2007 and December 2015, while for analysis (3) they are selected from  
 52 the 11-year point-source sample that includes events until 2017 [21]. The samples include neutrino  
 53 charged- and neutral-current interactions of all flavors. At energies of 10 TeV, the median angular  
 54 resolution for muon neutrinos is below  $0.5^\circ$ . In particular, analyses (1) and (2) require an event  
 55 signalness  $> 40\%$ , where the sig-  
 56 nalness is defined as the ratio  
 57 of the number of expected as-  
 58 trophysical events over the sum  
 59 of the expected atmospheric and  
 60 astrophysical events at a given  
 61 energy proxy, where a spec-  
 62 trum  $\phi = 1.01(E/100\text{TeV})^{-2.19}$ .  
 63  $10^{-18}\text{GeV}^{-1}\text{cm}^{-2}\text{sr}^{-1}$  was used  
 64 [22]. This selection results in a to-  
 65 tal of three tracks and no cascades.

66 **The Pierre Auger Observa-**  
 67 **tory** [23] is located in Argentina at  
 68 an average latitude of  $\sim 35.2^\circ$  and  
 69 a mean altitude of  $\sim 1400$  m above  
 70 the sea level. The Observatory is a  
 71 hybrid detector combining the in-  
 72 formation from a large surface detector array (SD) and a fluorescence detector (FD). The SD array,  
 73 spread over an area of  $3000$  km<sup>2</sup>, is composed of 1660 water-Cherenkov detectors. The FD array  
 74 consists of 27 telescopes at five peripheral buildings viewing the atmosphere over the SD array.  
 75 The data sample used in this work consists of 324 events observed with the SDs from January 2004  
 76 to April 2017 with reconstructed energies  $> 52$  EeV and zenith angle  $\theta \leq 80^\circ$  [24], which trans-  
 77 lates into a field of view ranging from  $-90^\circ$  to  $+45^\circ$  in declination. At these energies the angular  
 78 uncertainty is less than  $0.9^\circ$  [25], the statistical uncertainty in the energy determination is better  
 79 than  $12\%$  [26] and the systematic uncertainty in the absolute energy scale is  $14\%$  [27].

80 **The Telescope Array** (TA) experiment [28], located in Utah (USA), detects cosmic rays with  
 81  $E > 10^{18}$  eV. The surface array, composed by more than 500 scintillator detectors, extends over  
 82  $700$  km<sup>2</sup> of desert. In addition, there are three fluorescence telescope stations, instrumented with  
 83 12-14 telescopes each. The exposure of the detector covers the Northern Hemisphere and the South-  
 84 ern Hemisphere up to  $-15^\circ$ . A total of 143 events with energy  $\geq 57$  EeV and zenith angle  $\leq 55^\circ$ ,  
 85 recorded from May 2008 to May 2017, are used in this work [29]. These events have about  $1.5^\circ$   
 86 angular resolution,  $\sim 20\%$  energy resolution and a  $\sim 22\%$  systematic uncertainty on the energy



**Figure 1:** The UHECR events from TA and Auger are shown as orange and blue dots, respectively. The neutrino track- and cascade-like events from IceCube (HESE [10], EHE [11], 7-year through-going muons [12] samples) and ANTARES [20] are shown as black empty diamonds and crosses, respectively.

87 scale [29].

88 To account for the systematic energy shift in the absolute energy scale of UHECRs at the  
 89 energies of interest in this work, in the likelihood analyses presented here, the Auger energy scale  
 90 has been shifted by +14% and the TA energy scale by -14% following the latest studies of the  
 91 Auger-TA joint working group [30].

### 92 3. Analysis methods

93 In the following, the three analyses are described separately. In general, shower- and track-like  
 94 events are considered separately due to their different angular resolutions. Hence, separate  $p$ -values  
 95 are provided by each of the analyses.

#### 96 3.1 Cross-correlation analysis

97 The cross-correlation analysis counts the number,  $n_{\text{obs}}$ , of UHECR-neutrino pairs separated by  
 98 less than an angular distance,  $\Delta\alpha$ . This number is compared to the simulated number,  $n_{\text{exp}}$ , of pairs  
 99 within the same  $\Delta\alpha$  distance which are expected in the null-hypothesis scenario. Two separate  
 100 null-hypotheses are investigated: (i) an isotropic distribution of UHECRs, obtained by generating  
 101 isotropic CR datasets following the exposure of the two experiments and (ii) an isotropic distri-  
 102 bution of neutrinos, obtained by assigning randomly generated right-ascension values to the real  
 103 neutrino events, hence preserving the declination-dependent acceptance of the neutrino observato-  
 104 ries. The analysis is performed for different  $\Delta\alpha$  values, from  $1^\circ$  to  $30^\circ$  in  $1^\circ$  steps. The fraction  
 105 of isotropic simulations with equal or larger number of pairs than in data gives a measurement of  
 106 the probability (local  $p$ -value) that an observed excess of events arises by chance from an isotropic  
 107 distribution. The final  $p$ -value of the most important excess is evaluated by accounting for the  
 108 scan in angle. This analysis does not require any assumption on the (Galactic) magnetic field (un-  
 109 like the two analyses discussed in the following) since the scan on the angular distances between  
 110 the neutrinos and the CRs already accounts for any possible angular separation due to magnetic  
 111 deflections.

#### 112 3.2 Neutrino-stacking correlation analysis with UHECR directions

This analysis performs an unbinned-likelihood method by stacking the arrival directions of  
 the neutrinos and searching for coincident sources of cosmic rays (CRs). The signal hypothesis  
 assumes that UHECRs are correlated with high-energy neutrino directions. The background hy-  
 pothesis is consistent with an isotropic distribution of UHECRs across the sky. The logarithm of  
 the likelihood function is defined as:

$$\ln \mathcal{L}(n_s) = \sum_{i=1}^{N_{\text{Auger}}} \ln \left( \frac{n_s}{N_{\text{CR}}} S_{\text{Auger}}^i + \frac{N_{\text{CR}} - n_s}{N_{\text{CR}}} B_{\text{Auger}}^i \right) + \sum_{i=1}^{N_{\text{TA}}} \ln \left( \frac{n_s}{N_{\text{CR}}} S_{\text{TA}}^i + \frac{N_{\text{CR}} - n_s}{N_{\text{CR}}} B_{\text{TA}}^i \right), \quad (3.1)$$

where  $n_s$  is the number of signal events, i.e. UHECRs correlated with neutrino directions, and  
 is the only free parameter;  $N_{\text{CR}} = N_{\text{Auger}} + N_{\text{TA}}$  is the total number of CR events.  $S^i$  and  $B^i$  are,  
 respectively, the signal and background probability distribution functions (PDFs) for each CR ob-  
 servatory. The signal PDF, for the  $i^{\text{th}}$  CR at a given direction  $\vec{r}_i$  and with energy  $E_i$ , can be expressed

as:

$$S_{\text{CR observatory}}^i(\vec{r}_i, E_i) = R_{\text{CR observatory}}(\delta_i) \sum_{j=1}^{N_{\text{src}}} S_j(\vec{r}_i, \sigma(E_i)), \quad (3.2)$$

113 where  $R_{\text{CR observatory}}$  is the relative experiment exposure at a given event declination,  $\delta_i$ ,  $N_{\text{src}}$  is the  
 114 number of stacked (neutrino) sources and  $S_j(\vec{r}_i, \sigma(E_i))$  is the value of the normalized directional  
 115 likelihood map for the  $j^{\text{th}}$  source taken at position  $\vec{r}_i$ . The arrival direction of the  $i^{\text{th}}$  UHECR event is  
 116 obtained by smearing the source position with a two-dimensional Gaussian function with standard  
 117 deviation  $\sigma(E_i)$  calculated as  $\sigma(E_i) = \sqrt{\sigma_{\text{CR observatory}}^2 + \sigma_{\text{MD}}^2}$ , where  $\sigma_{\text{CR observatory}}$  is the angular  
 118 resolution of the CR observatory ( $0.9^\circ$  for Auger and  $1.5^\circ$  for TA) and  $\sigma_{\text{MD}} = D \times 100 \text{ EeV}/E_{\text{CR}}$  is  
 119 the energy-dependent Galactic magnetic deflection. In the analysis,  $D$  is assigned three benchmark  
 120 values. In order to account for the differences of the Galactic magnetic field in the Northern and  
 121 Southern hemispheres, two average deflection values are calculated, whereas previous analyses  
 122 used all-sky average deflection values [31, 32], by considering the Galactic magnetic field models  
 123 of Pshirkov et al. [33] and Jansson and Farrar [34]. Assuming a pure proton-like CR sample with  
 124  $E_{\text{CR}} = 100 \text{ EeV}$ , mean angular deflection values of  $2.4^\circ$  and  $3.7^\circ$  are obtained, for the North and  
 125 South respectively. To account for possible heavier compositions of the CRs or larger contributions  
 126 of the magnetic fields, the average deflection values in the North and in the South at  $E_{\text{CR}} = 100$   
 127 EeV are increased by factors 2 and 3. Finally, the background PDFs,  $B_{\text{Auger}}$  and  $B_{\text{TA}}$  in Eq. 3.1,  
 128 represent the probability of observing a cosmic ray from a given direction assuming an isotropic  
 129 flux. Therefore they are calculated from the Auger and TA normalized exposures. The test statistic  
 130 (TS) is defined as  $\text{TS} = 2 \ln (\mathcal{L}(\hat{n}_s)/\mathcal{L}(\hat{n}_s = 0))$ , where  $\hat{n}_s$  denotes the optimized parameter.

### 131 3.3 UHECR-stacking correlation analysis with neutrino directions

This analysis is based on an unbinned likelihood method for searching point-like neutrino sources [17], with additional information from stacking UHECR arrival directions. The neutrino events are weighted according to the relative experiment exposure [35]. The signal hypothesis assumes point-like neutrino sources to be spatially correlated with UHECR arrival directions, which are subject to a specific magnetic deflection hypothesis. The background hypothesis assumes that neutrino events are uniformly distributed over the whole sky. The free signal parameters of the likelihood function,  $\mathcal{L}$ , are the numbers of neutrino signal events,  $n_s$ , and the spectral indices,  $\gamma_s$ , for each possible neutrino source at positions  $\vec{x}_s$ . The logarithm of the likelihood function is defined as:

$$\ln \mathcal{L} = \underbrace{\sum_{s=1}^{N_{\text{CR}}}}_{\text{stacking, step 3}} \left[ \underbrace{\left( \sum_{i=1}^{N_{\text{V}}} \ln \left( \frac{n_s}{N_{\text{V}}} S_i(\gamma_s, \vec{x}_s) + \left( 1 - \frac{n_s}{N_{\text{V}}} \right) B_i(\vec{x}_s) \right)}_{\text{neutrino data, step 1}} - \underbrace{\frac{(\vec{x}_s - \vec{x}_{\text{CR},s})^2}{2\sigma(E_{\text{CR},s})^2}}_{\text{UHECR data, step 2}} \right]. \quad (3.3)$$

132 The first part of the likelihood formula (step 1 in eq. 3.3) is determined by information from neu-  
 133 trino data, where the sum runs over all experimentally measured neutrino candidates  $N_{\text{V}}$ . The sig-  
 134 nal PDF,  $S_i$ , describes a point-like neutrino source at position  $\vec{x}_s$  with  $n_s$  events following a certain  
 135 spectral index,  $\gamma_s$ . The index  $s$  denotes one neutrino source as counterpart to one UHECR event, as  
 136 described later. The background PDF,  $B_i$ , is determined from experimental neutrino events whose

137 right ascension coordinates were assigned random values. Any information from UHECR data is  
 138 contained in the spatial prior functions shown in the second part of the likelihood (step 2). In the  
 139 first step, the signal parameters  $(n_s, \gamma_s)$  are optimized *without the spatial prior function* on grid  
 140 positions covering the whole sky, based on IceCube’s standard procedure for searching for point-  
 141 like sources [17]. The result is translated into a TS map of the neutrino sky. The TS is defined  
 142 as  $\text{TS}(\vec{x}_s) = 2 \ln(\mathcal{L}_{\text{step 1}}(\hat{n}_s, \hat{\gamma}_s) / \mathcal{L}_{\text{step 1}}(n_s = 0))$ , where  $(\hat{n}_s, \hat{\gamma}_s)$  denote the optimized parameters.  
 143 In a second step, the arrival direction of one UHECR event,  $\vec{x}_{\text{CR},s}$ , and the corresponding smear-  
 144 ing,  $\sigma(E_{\text{CR},s})$  (cf. sec. 3.2) are used to construct a 2D-Gaussian function, which is logarithmically  
 145 added to the TS map. This results in an effective selection of the neutrino sky where the largest re-  
 146 maining TS spot is the most likely neutrino source counterpart to the selected UHECR event. This  
 147 is equivalent to optimizing  $n_s, \gamma_s$  and the source position in presence of a spatial Gaussian prior  
 148 function. The third step is to repeat the procedure for all selected UHECRs, and the resulting TS  
 149 values are summed to yield the final TS. This is equivalent to a stacking of independent neutrino  
 150 sources selected by UHECR and relative deflection information. Note that this procedure ensures  
 151 that each UHECR event has one neutrino-source counterpart in its vicinity, while allowing one neu-  
 152 trino source to be counterpart to several UHECRs in its vicinity. Three different lower energy cuts  
 153  $E_{\text{CR}} \geq [70, 85, 100]$  EeV are applied to the combined UHECR sample in order to study a potential  
 154 energy dependence in the final TS. Magnetic deflection values of  $D = 3^\circ$  and  $6^\circ$  are used uniformly  
 155 over the whole sky.

## 156 4. Results

### 157 4.1 Cross-correlation analysis

158 The results of the scan in angle are shown in Fig. 2, where the relative number of observed  
 159 pairs with respect to the expected value from an isotropic distribution of neutrinos is shown for  
 160 tracks (left) and cascades (right). The maximum departure from the isotropy is found at  $14^\circ$  for  
 161 tracks, where 582 pairs are observed, and at  $16^\circ$  for cascades, with observed 763 pairs. The post-  
 162 trial  $p$ -values are 0.23 for tracks and 0.15 for cascades. The maximum departure from isotropy  
 163 with respect to an isotropic distribution of UHECRs is found at  $10^\circ$  for tracks, where 303 pairs  
 164 are observed and at  $16^\circ$  for cascades, with observed 763 pairs. The post-trial  $p$ -values are 0.84 for  
 165 tracks and 0.18 for cascades.

### 166 4.2 Neutrino-stacking correlation analysis with UHECR directions

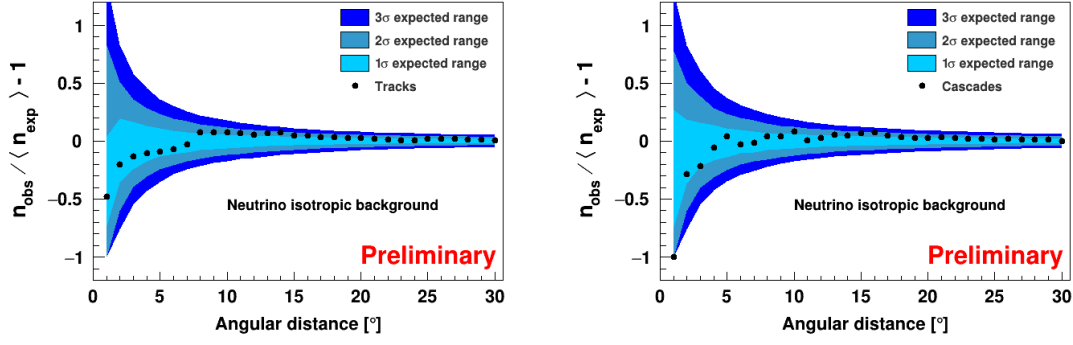
167 The results are shown in Tab. 1. The most significant deviation from an isotropic flux of CRs  
 168 occurs for the magnetic deflection parameter set of  $D = (7.2^\circ, 11.1^\circ)$  with the high-energy cascade  
 events. The observed pre-trial  $p$ -value of 0.29 corresponds to 0.90 post-trial, obtained by consider-

$D$	$(2.4^\circ, 3.7^\circ)$	$(4.8^\circ, 7.4^\circ)$	$(7.2^\circ, 11.1^\circ)$
tracks	underfluctuation	underfluctuation	underfluctuation
cascades	underfluctuation	0.41	0.29

**Table 1:** Pre-trial  $p$ -values for the neutrino-stacking analysis with the samples of high-energy tracks and cascades assuming an isotropic flux of UHECRs.

169 ing multiple realizations of pseudo-experiments of randomly distributed CRs with  $D = (7.2^\circ, 11.1^\circ)$   
 170





**Figure 2:** Relative excess of pairs,  $n_{\text{obs}}/\langle n_{\text{exp}} \rangle - 1$ , as a function of the maximum angular separation between the neutrino and UHECR pairs, for track- (left) cascade-like (right) events in the case of an isotropic distribution of neutrinos. The different color bands stand for the regions containing the 1, 2 and  $3\sigma$  fluctuations from an isotropic distribution.

171 and by accounting for the trial factor due to having tested three different sets of magnetic de-  
 172 flections. Given these numbers, no sign of correlations in the arrival directions of UHECRs and  
 173 neutrinos is found.

#### 174 4.3 UHECR-stacking correlation analysis with neutrino directions

175 Six  $p$ -values for each of the signal hypotheses described in sec. 3.3 are calculated with respect  
 176 to an isotropic neutrino flux, summarized in Tab. 2. All six  $p$ -values are well-compatible with the  
 177 background expectation. The smallest  $p$ -value (6% for  $D = 6^\circ$  and  $E_{\text{CR}} \geq 85$  EeV), after correction  
 for the six correlated tests, becomes 16%.

$D$ [ $^\circ$ ]	3			6		
$E_{\text{CR}}$ [EeV] $\geq$	70	85	100	70	85	100
$p$ -value	0.27	0.46	0.84	0.10	0.06	0.39

**Table 2:** All pre-trial  $p$ -values for different UHECR energy cuts and deflection hypotheses.

178

## 179 5. Discussion

180 The results of the three analyses presented here do not allow to conclude about the presence of  
 181 possible correlations between arrival directions of UHECRs and high-energy neutrinos. Previous  
 182 analyses on reduced data sets had shown a post-trial  $p$ -value of  $5.0 \times 10^{-4}$  for the cascades with a  
 183 cross-correlation analysis, under the assumption of an isotropic flux of UHECRs, and of  $8.0 \times 10^{-4}$   
 184 for cascades with the neutrino-stacking analysis under a deflection hypothesis of  $D = 6^\circ$  [31]. The  
 185 absence of correlation found with the current data samples and discussed analysis hypotheses must  
 186 be carefully interpreted, since it does not imply an absolute lack of correlation in the origin of  
 187 the two messengers. The main uncertainties in the current analyses are the poor knowledge of  
 188 the Galactic magnetic field and the not yet conclusive understanding of the CR composition. Fur-  
 189 thermore, due to the GZK horizon, the largest distances covered by UHECRs are not expected to  
 190 exceed 10-100 Mpc, depending on the CR composition. On the other hand, neutrinos can reach  
 191 us from cosmological distances. Finally, neutrinos originating at the cosmic rays acceleration sites

192 are expected to carry few percents of the energy of the original cosmic ray. Thus, the neutrinos ob-  
 193 served by IceCube and ANTARES might have been produced by cosmic rays of much lower energy  
 194 than the ones in the datasets by Auger and TA. For these reasons, only a few percent of neutrinos  
 195 may be expected to originate from the same astrophysical sources of the detected UHECRs.

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